

Observation of electroweak Z+2j production at ATLAS

[a.k.a Z-boson production via weak boson fusion]

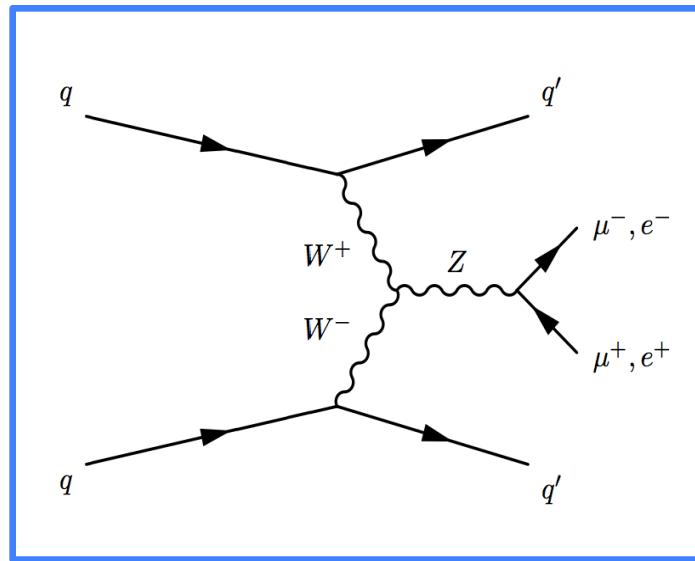
Andrew Pilkington – University College London

Presented at the workshop on multi-boson interactions, BNL, October 2014

Outline

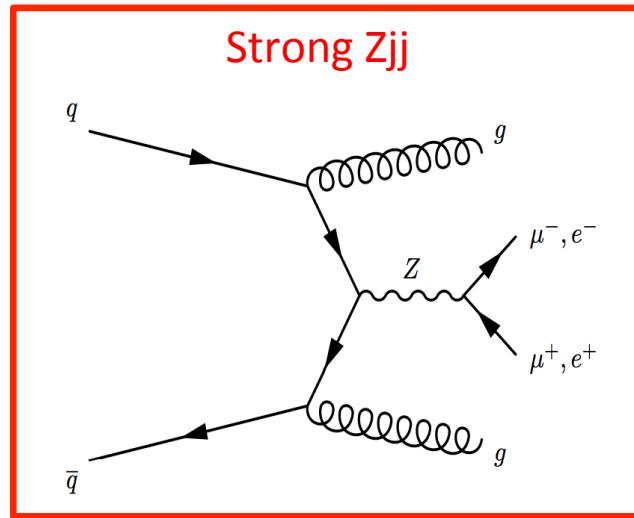
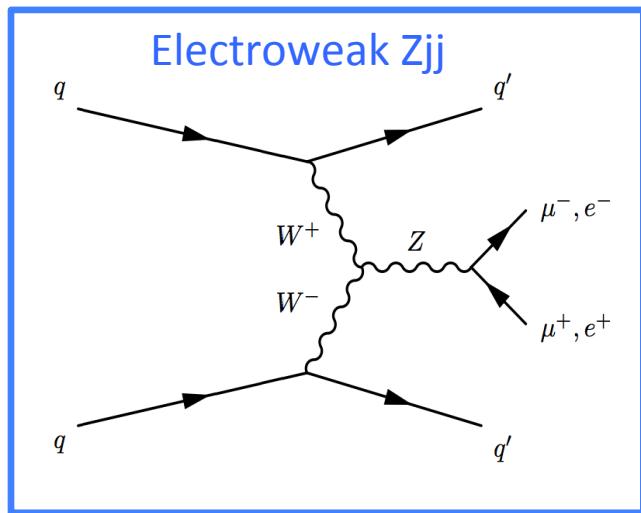
- 1) Cross sections and distributions for inclusive Z+2j production
- 2) Extracting the electroweak Z+2j component

Why bother with electroweak Z+2j production?



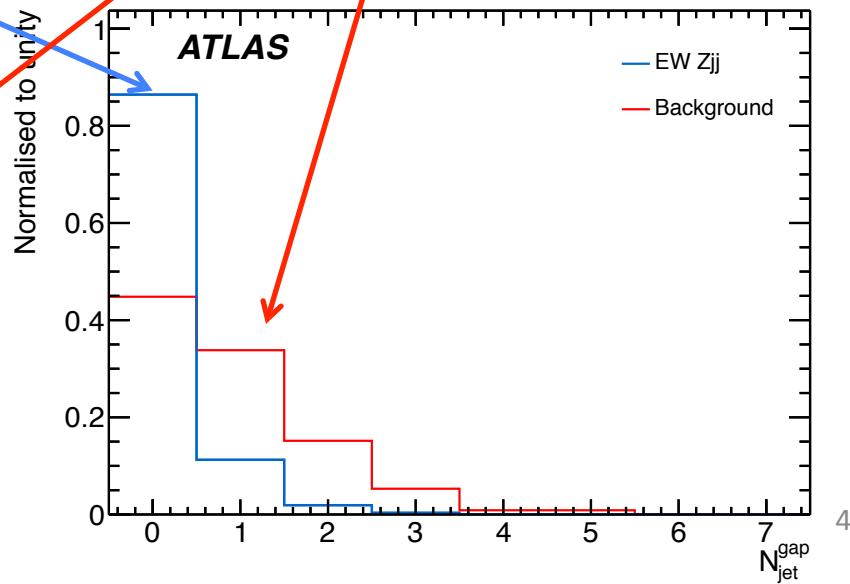
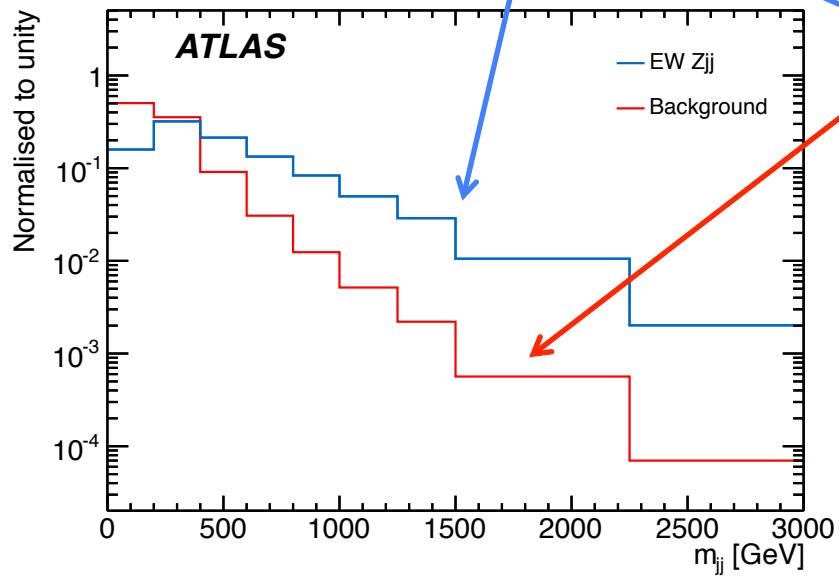
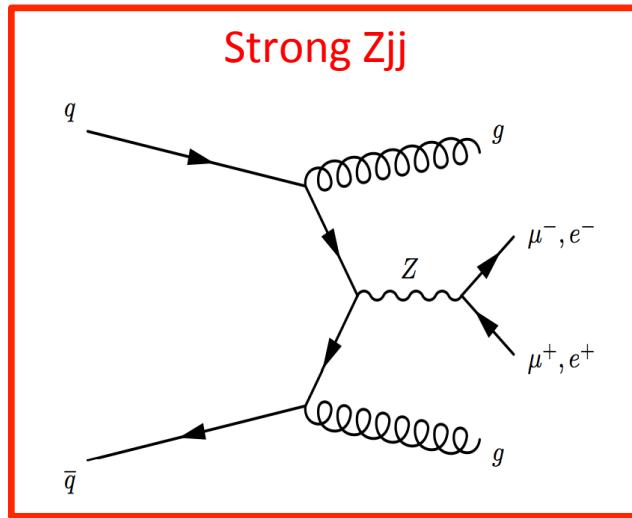
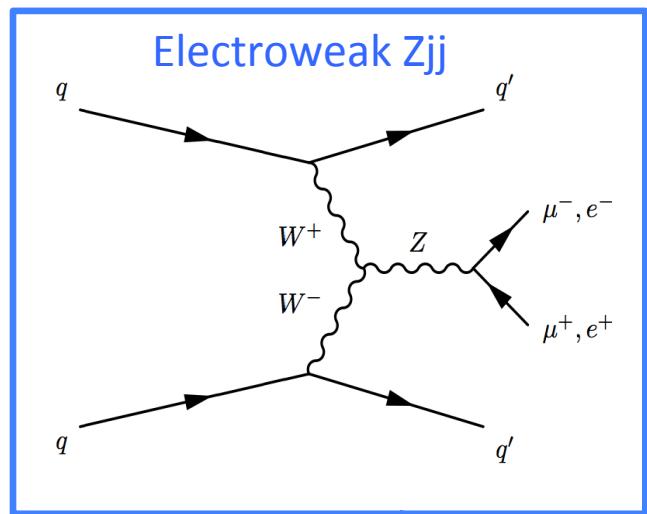
- Weak-boson fusion/scattering plays an important role in Higgs measurements (couplings, $\pi\pi$) and searches for anomalous quartic gauge couplings.
- Z-production via weak boson fusion is a *standard candle* for these processes
 - first observation of weak boson fusion at a hadron collider
 - a direct test of the ZWW coupling.

Z+2j production at the LHC



- Electroweak Z+2j production is *rare*: only $\sim 1\%$ of the inclusive Z+2j cross section
- Electroweak Z+2j has two characteristic features:
 - Dijet system covering a large rapidity interval and large invariant mass
 - Little additional jet activity in the rapidity interval

Z+2j production at the LHC



Cross sections and distributions for inclusive Z+2j production

Detector-corrected measurements at $\sqrt{s}=8\text{TeV}$

- 1) Cross sections measured in five fiducial regions
- 2) Differential distributions sensitive to dijet kinematics (m_{jj} , Δy)
- 3) Differential distributions sensitive to ‘in-gap’ jet activity

Five fiducial regions: different sensitivity to electroweak Zjj

Object	<i>baseline</i>	<i>high-mass</i>	<i>search</i>	<i>control</i>	<i>high-p_T</i>
Leptons		$ \eta^\ell < 2.47, p_T^\ell > 25 \text{ GeV}$			
Dilepton pair		$81 \leq m_{\ell\ell} \leq 101 \text{ GeV}$			
	—		$p_T^{\ell\ell} > 20 \text{ GeV}$	—	
Jets		$ y^j < 4.4, \Delta R_{j,\ell} \geq 0.3$			
		$p_T^{j1} > 55 \text{ GeV}$			$p_T^{j1} > 85 \text{ GeV}$
		$p_T^{j2} > 45 \text{ GeV}$			$p_T^{j2} > 75 \text{ GeV}$
Dijet system	—	$m_{jj} > 1 \text{ TeV}$	$m_{jj} > 250 \text{ GeV}$	—	
Interval jets	—	—	$N_{\text{jet}} = 0$	$N_{\text{jet}} \geq 1$	—
Zjj system	—	—	$p_T^{\text{balance}} < 0.15$	$p_T^{\text{balance},3} < 0.15$	—

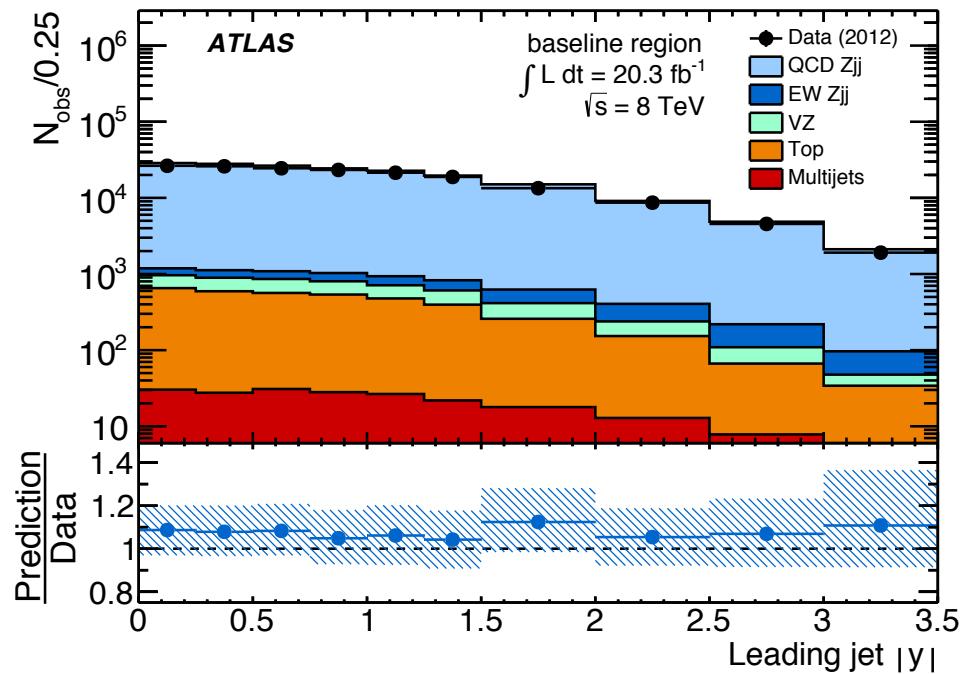
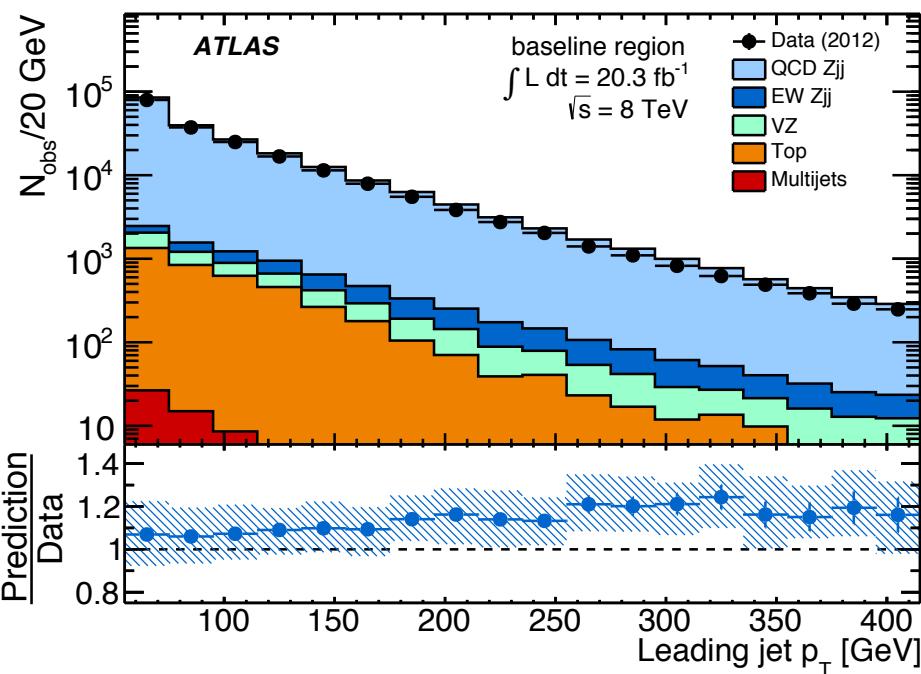
--- Z-boson selection

--- Baseline jet selection

--- Probe of high-p_T or high-mass

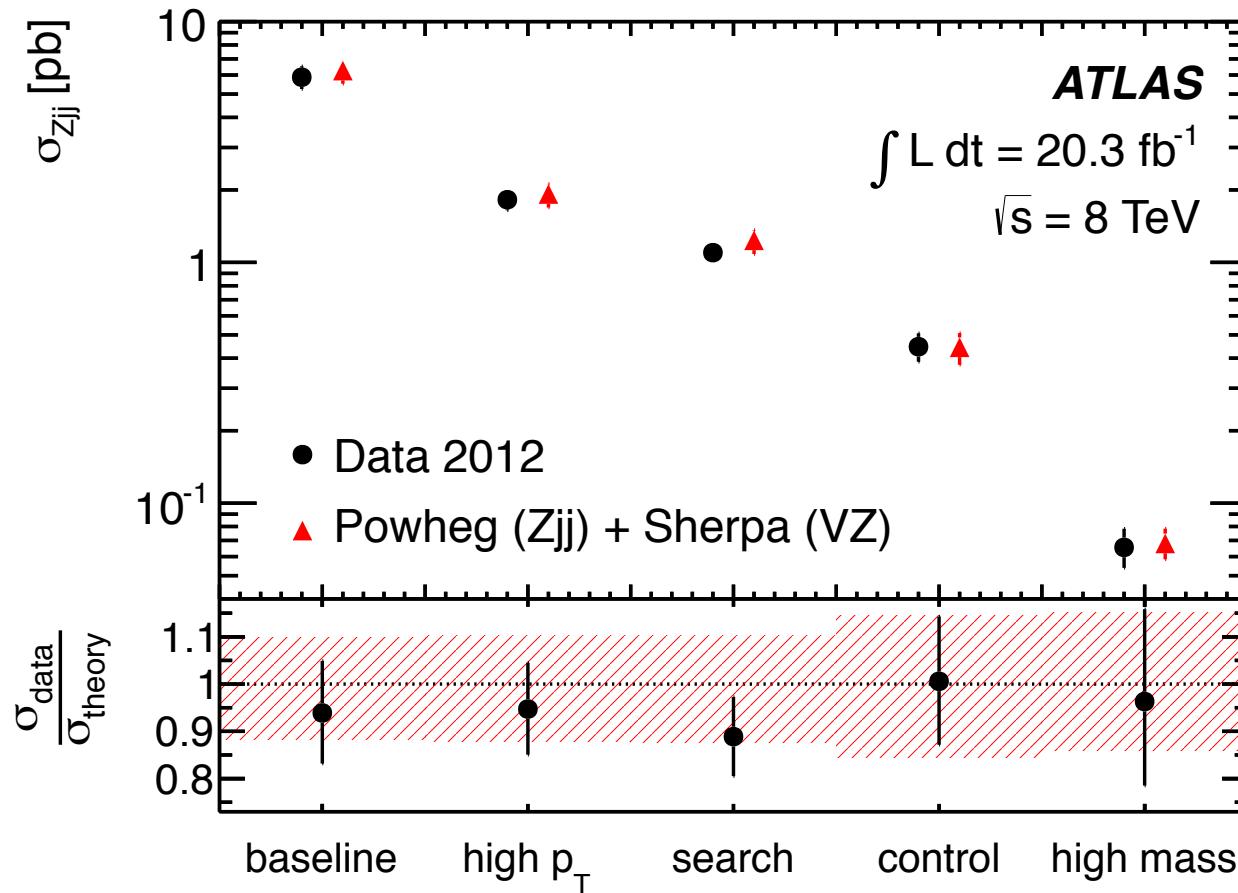
--- Search/control cuts for electroweak extraction

Example of data/simulation in the *baseline* region



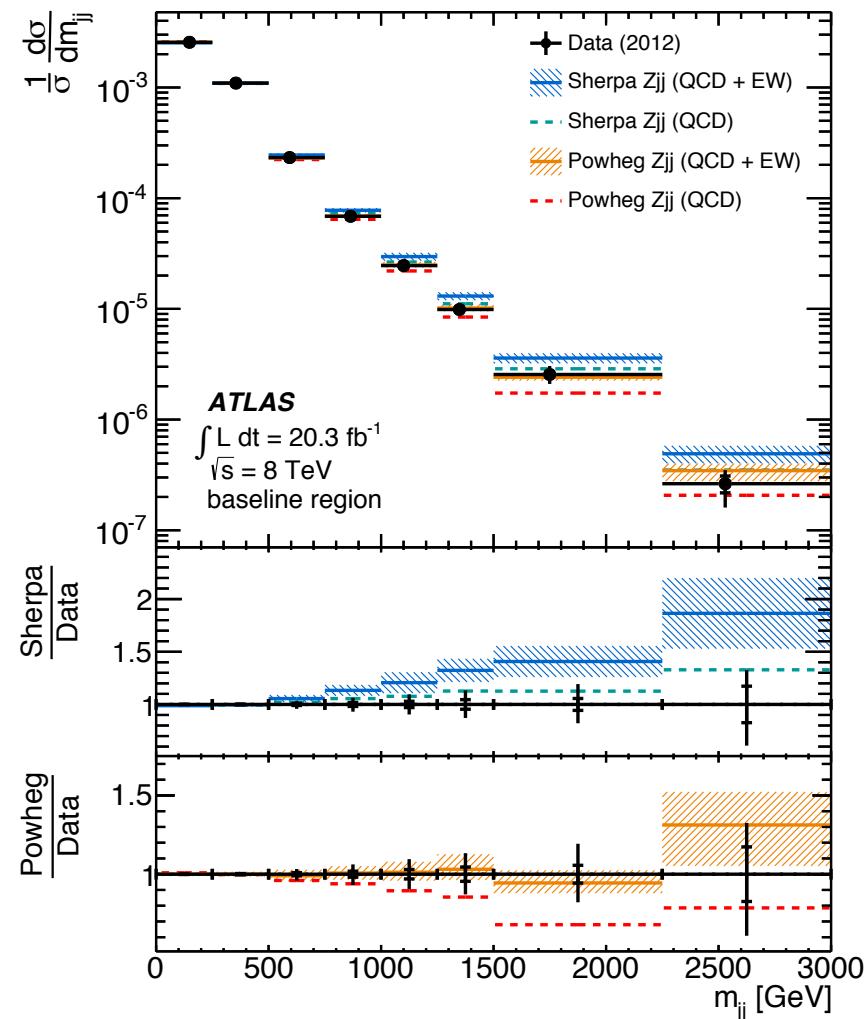
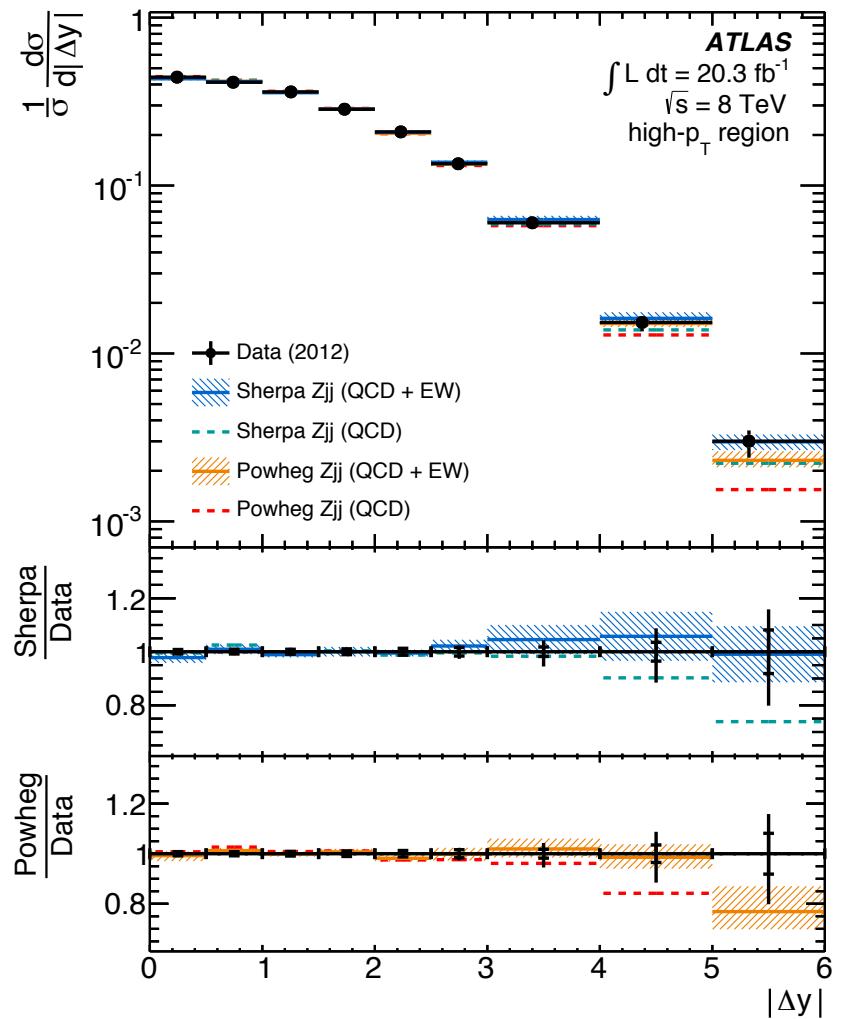
- Small non-Z backgrounds
- Reasonable agreement between data and simulation

Measurement of inclusive Z+2j cross sections (II)



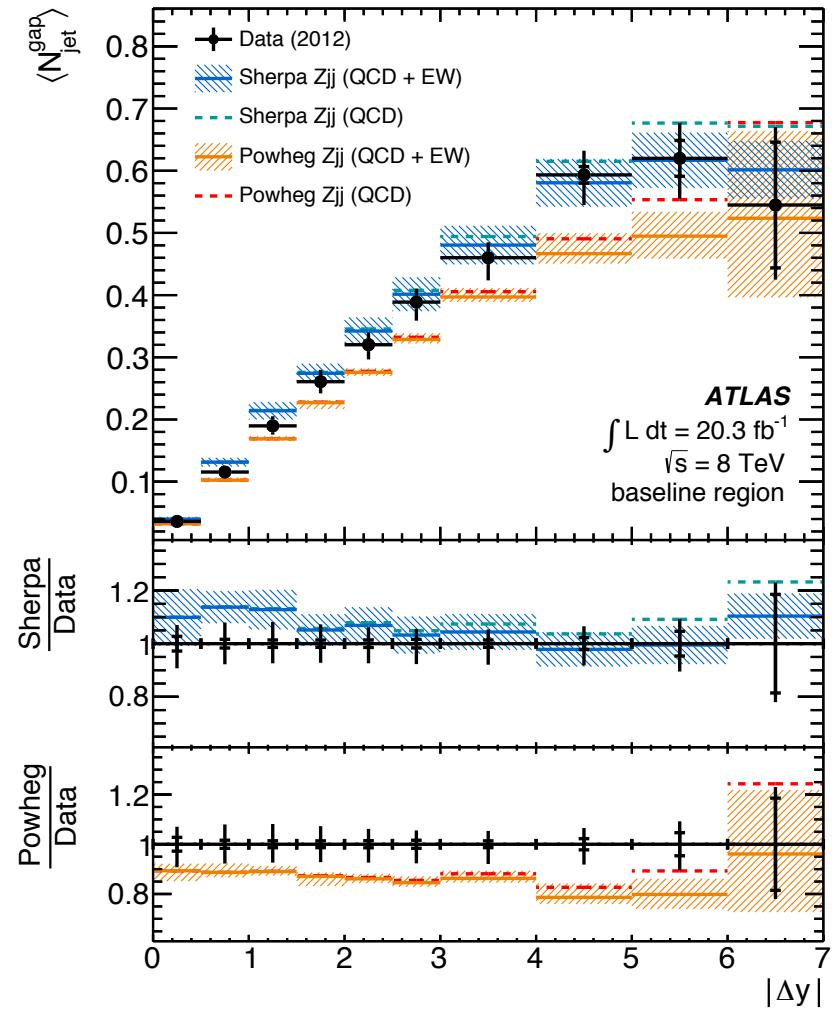
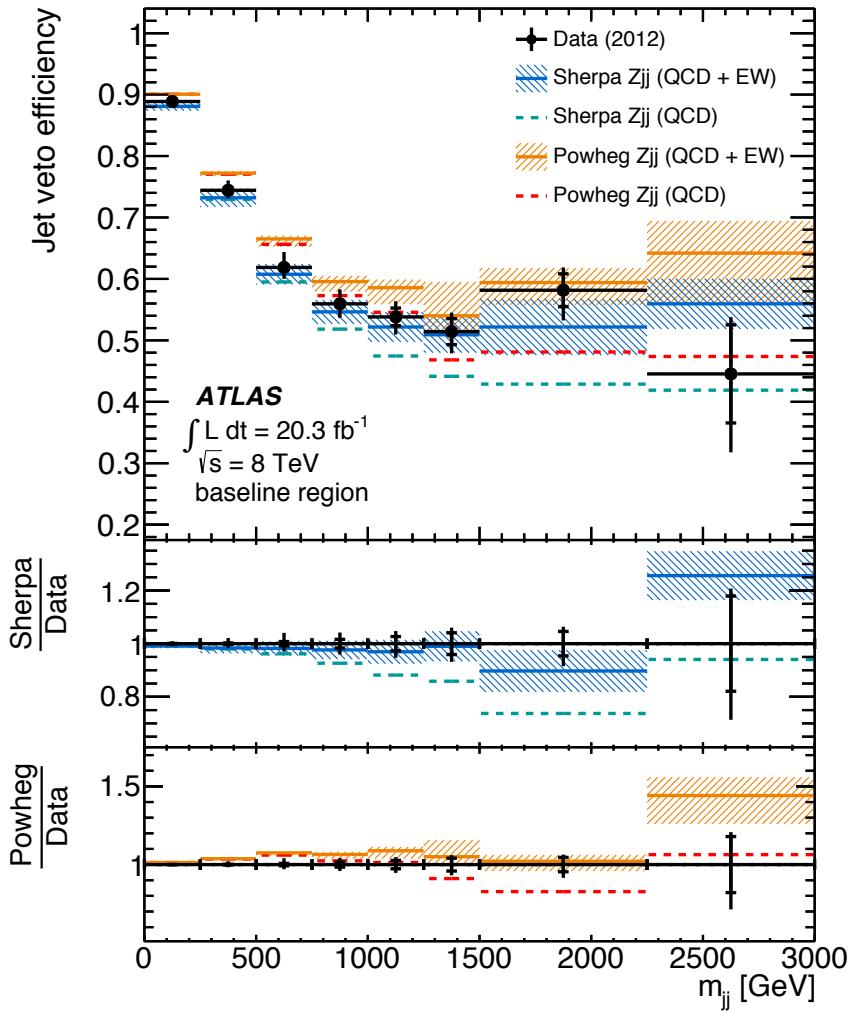
Measure cross section by $\sigma_{\text{fid}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\int L dt \cdot \mathcal{C}}$

Measurement of dijet kinematics (unfolded)



Powheg accurate to next-to-leading order (NLO) in QCD for Z+2j production
Sherpa accurate only to leading order (LO) in QCD for Z+2j production

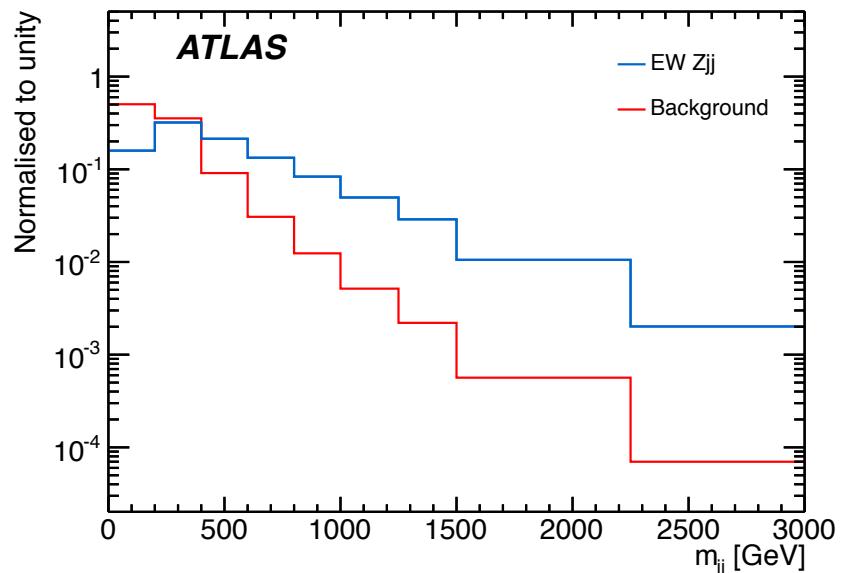
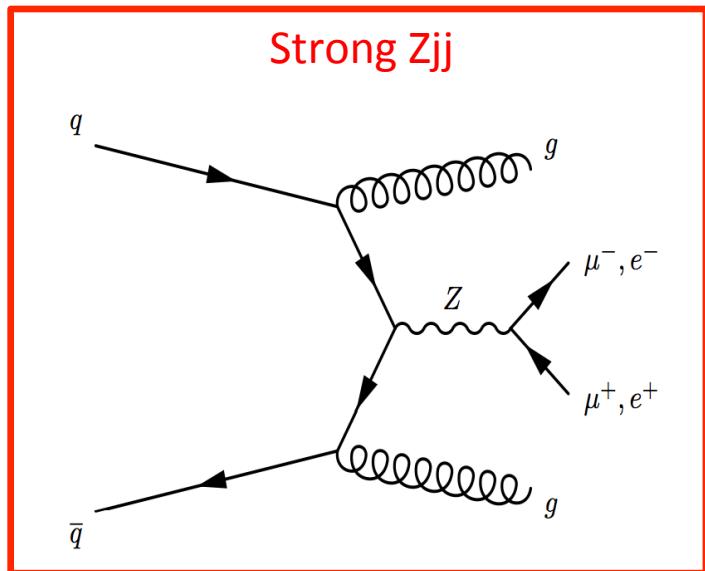
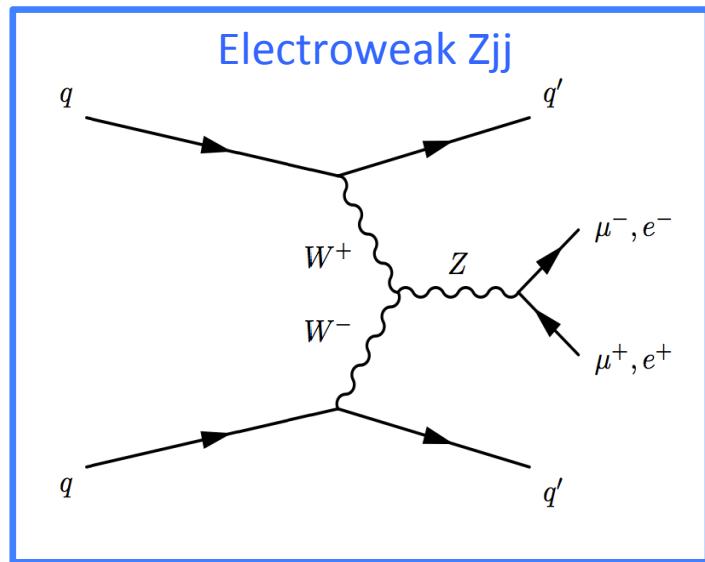
Measurement of in-gap jet activity (unfolded)



Sherpa accurate to LO in QCD for third and fourth jet emission
Powheg only accurate to LO in QCD for third jet emission

Extracting the electroweak Z+2j component ($\sqrt{s}=8\text{TeV}$)

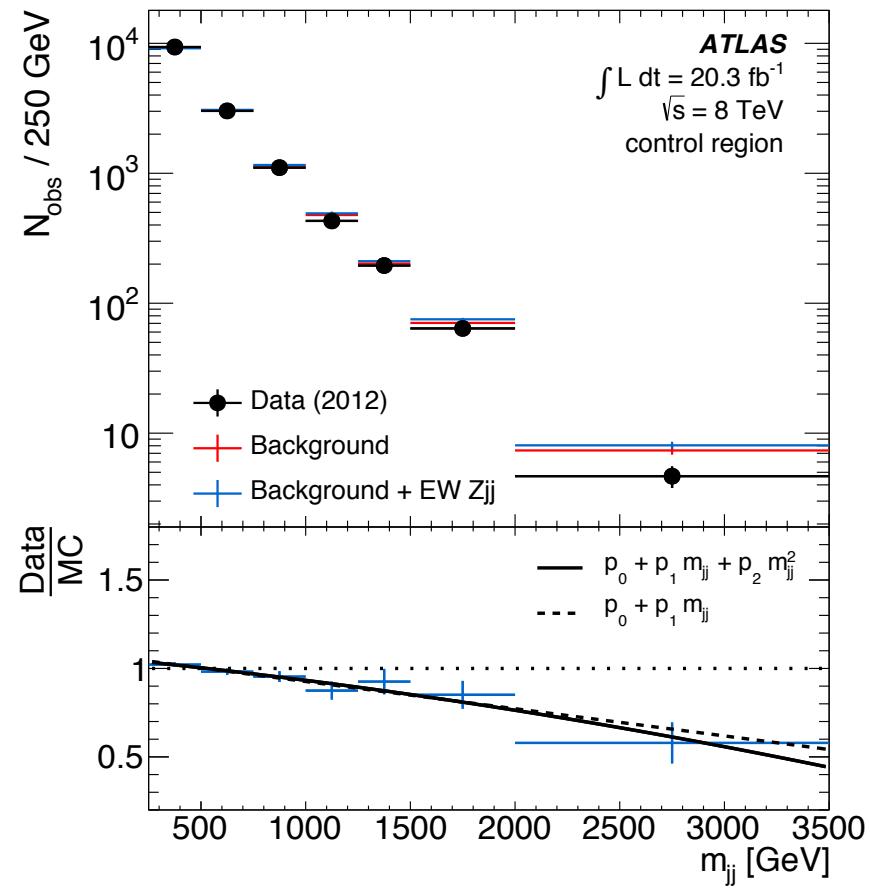
Extracting the signal - methodology



- Electroweak component is extracted by a two template fit to the dijet invariant mass
- This fit is carried out in the *search region*, which has a veto on additional central jet activity

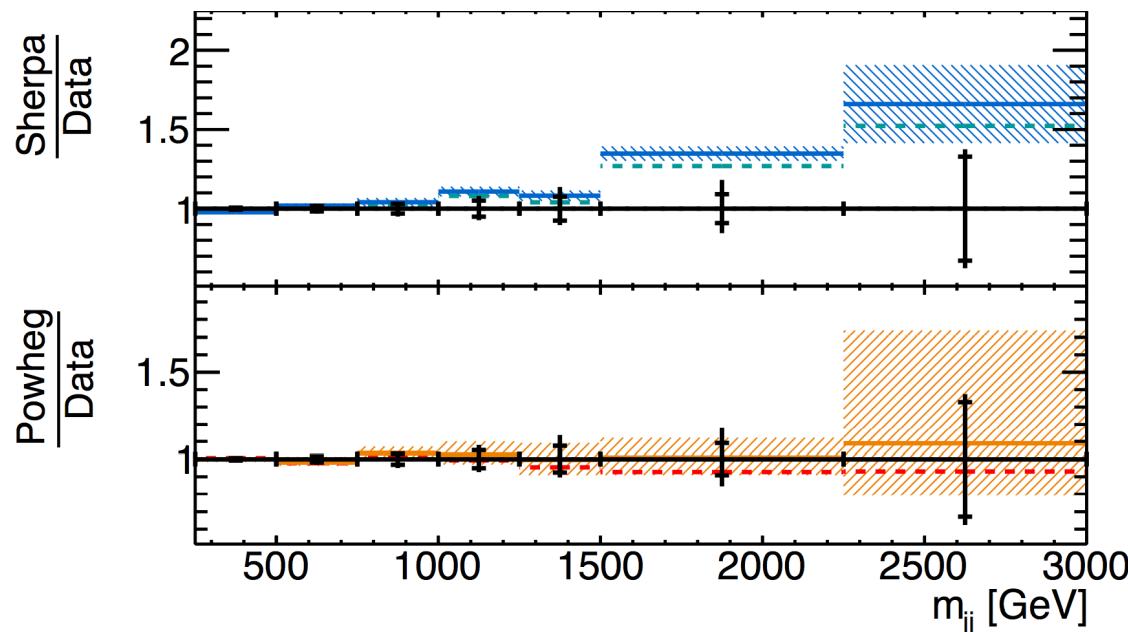
Data-driven constraint for the background model

- Control region defined by reversing the jet veto.
- Basic idea: correct the simulation in the *search* region using the data/MC ratio in the *control* region.
 - Plot shows correction derived for SHERPA
- Added bonus: limits the impact of jet energy scale uncertainties on the tagging jets
- Remaining experimental and theoretical uncertainties associated with the third jet modelling (extrapolation from *control* \rightarrow *search*)



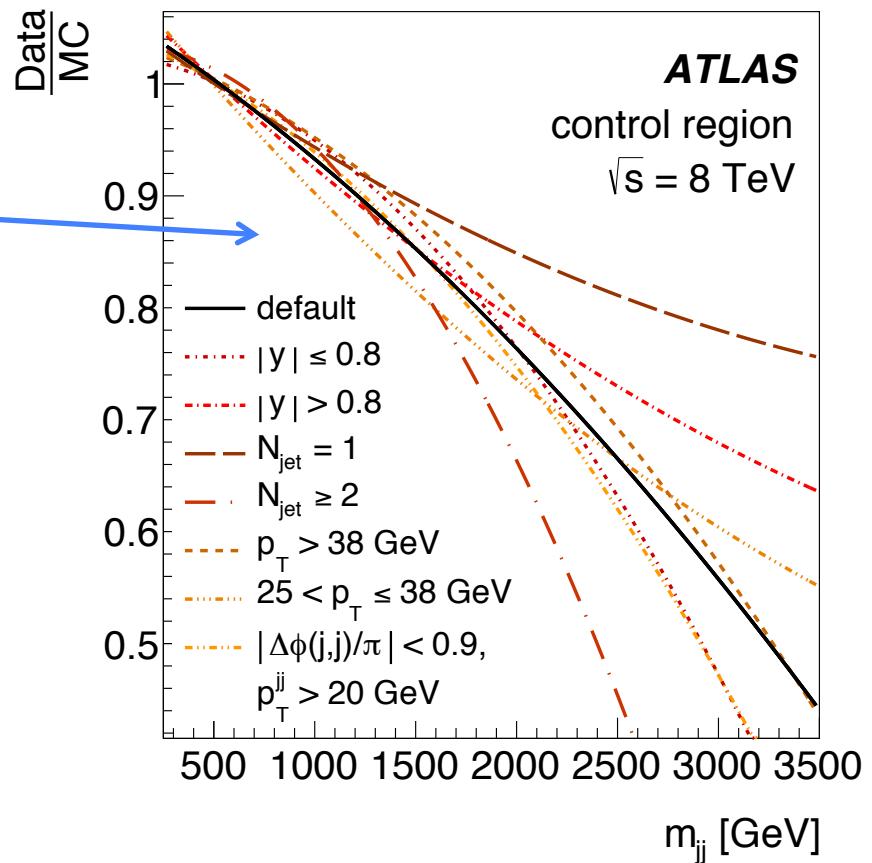
Data-driven constraint for the background model (I)

- Unfolded distributions show that POWHEG gives a better description than SHERPA as a function of m_{jj}
 - Choice of generator checked by reweighting SHERPA events to POWHEG in both the search and control regions.
 - Full analysis procedure repeated (i.e. new control region constraint derived)
 - Extracted signal yields agree to 0.8%



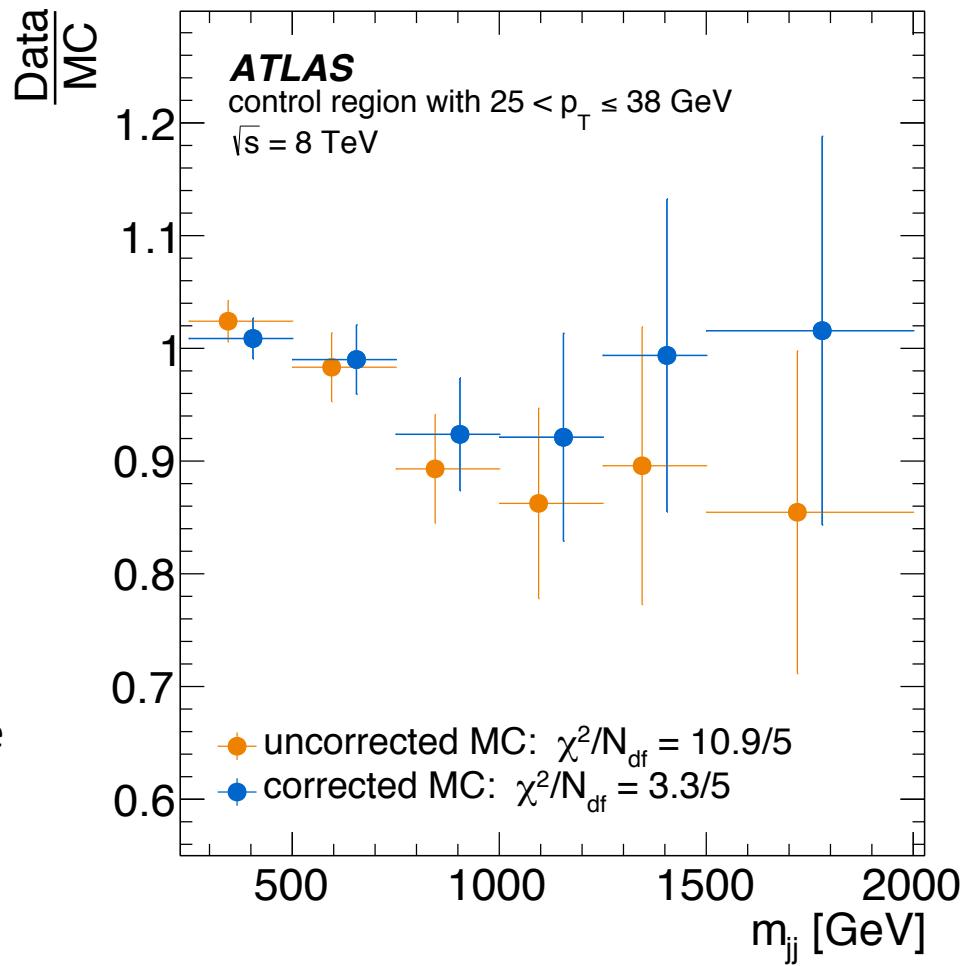
Data-driven constraint for the background model (II)

- Choice of control region validated by splitting it into seven sub-regions,
 - deriving new constraints,
 - repeating full analysis chain.
 - extracted signal yields agree to within 5%.



Data-driven constraint for the background model (III)

- The seven control regions are all signal-suppressed.
- Can use orthogonal sub-regions to test whether the reweighting function derived on one region improves the agreement between simulation and data in another region
- Agreement improved in all cases,
- Example, shown here is for orthogonal sub-regions that differ in the third jet transverse momentum.

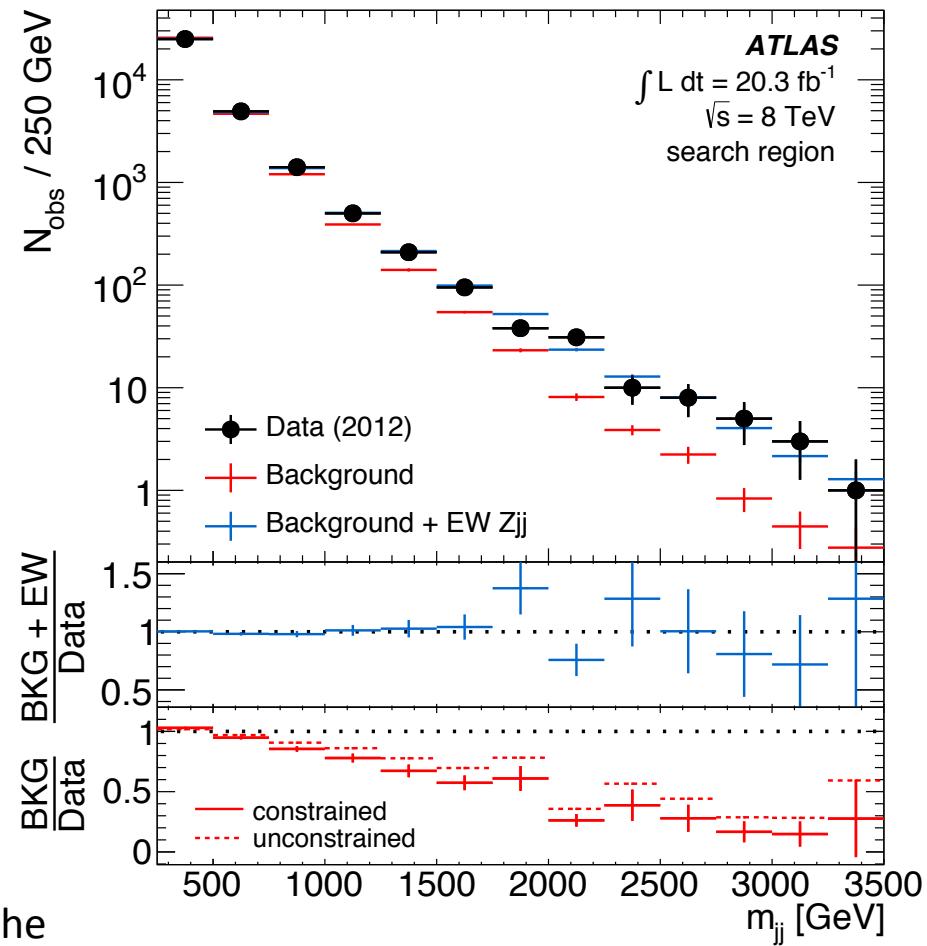


Extracting the signal - results

Fit results

Electron+muon	
Data	32186
MC predicted N_{bkg}	$32600 \pm 2600^{+3400}_{-4000}$
MC predicted N_{EW}	$1333 \pm 50 \pm 40$
Fitted N_{bkg}	$30530 \pm 216 \pm 40$
Fitted N_{EW}	$1657 \pm 134 \pm 40$

Background-only hypothesis rejected
at greater than 5σ significance



- Extracted yield converted to a cross section in the search fiducial region:

$$\sigma_{\text{EW}} = 54.7 \pm 4.6 \text{ (stat)} ^{+9.8}_{-10.4} \text{ (syst)} \pm 1.5 \text{ (lumi)} \text{ fb.}$$

Breakdown of the electroweak cross section systematics

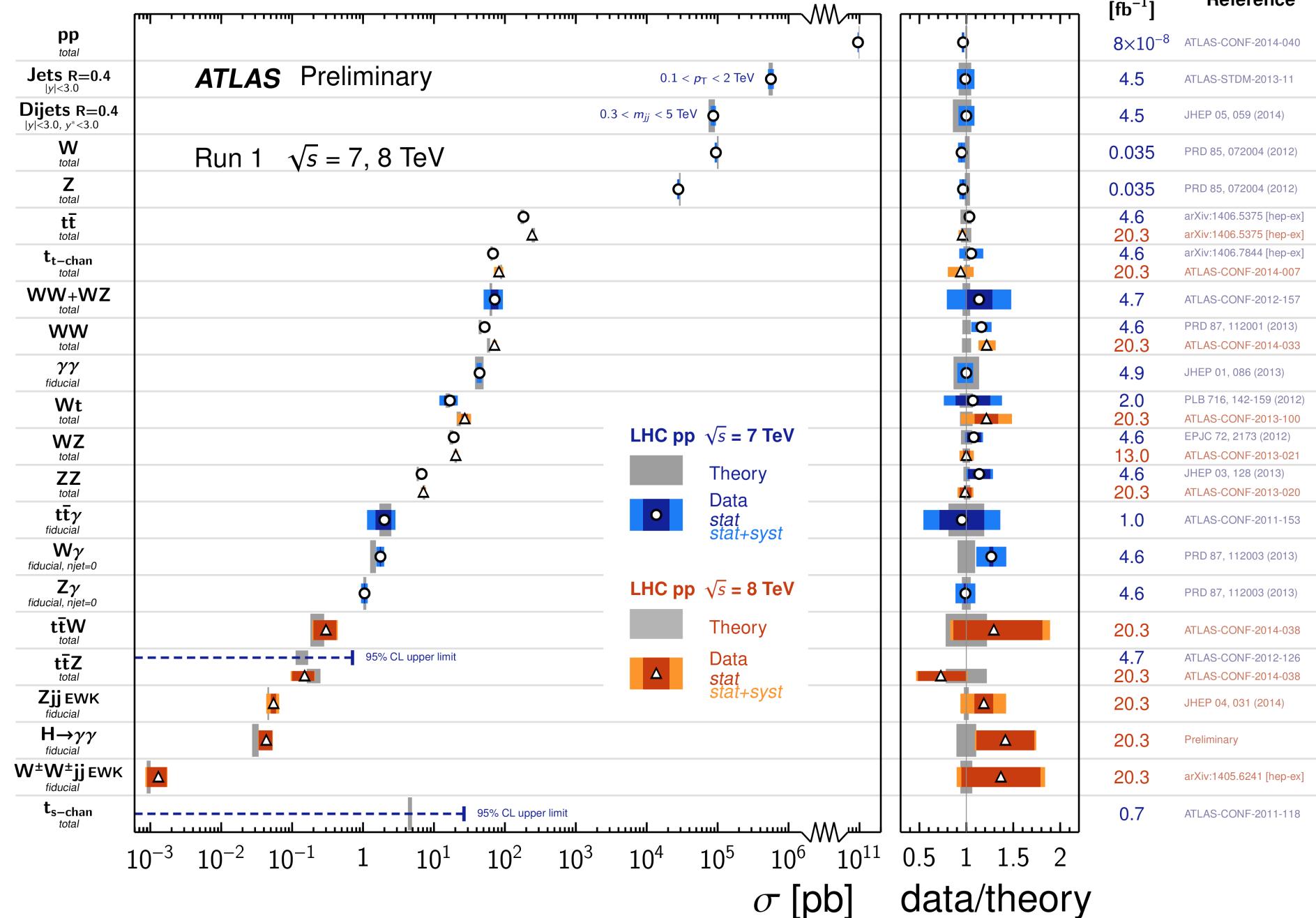
Source	ΔN_{EW}		$\Delta \mathcal{C}_{\text{EW}}$	
	Electrons	Muons	Electrons	Muons
Lepton systematics	—	—	$\pm 3.2 \%$	$\pm 2.5 \%$
Control region statistics	$\pm 8.9 \%$	$\pm 11.2 \%$	—	—
JES	$\pm 5.6 \%$		$\begin{array}{c} +2.7 \\ -3.4 \end{array} \%$	
JER	$\pm 0.4 \%$		$\pm 0.8 \%$	
Pileup jet modelling	$\pm 0.3 \%$		$\pm 0.3 \%$	
JVF	$\pm 1.1 \%$		$\begin{array}{c} +0.4 \\ -1.0 \end{array} \%$	
Signal modelling	$\pm 8.9 \%$		$\begin{array}{c} +0.6 \\ -1.0 \end{array} \%$	
Background modelling	$\pm 7.5 \%$		—	
Signal/background interference	$\pm 6.2 \%$		—	
PDF	$\begin{array}{c} +1.5 \\ -3.9 \end{array} \%$		$\pm 0.1 \%$	

Standard Model Production Cross Section Measurements

Status: July 2014

$\int \mathcal{L} dt$
[fb $^{-1}$]

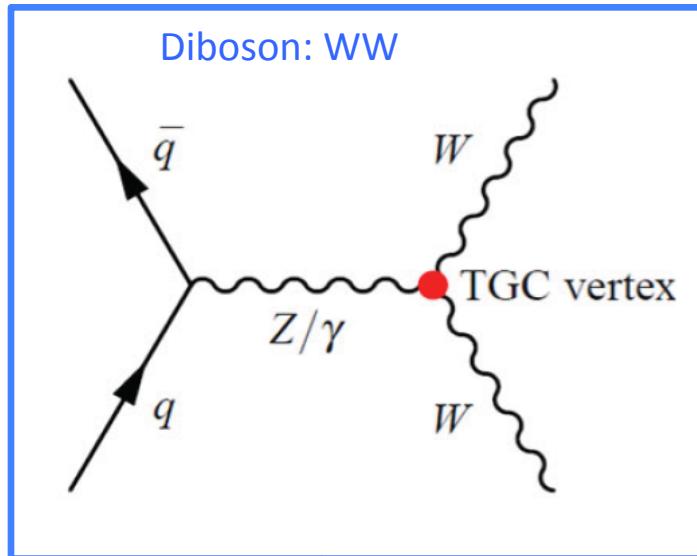
Reference



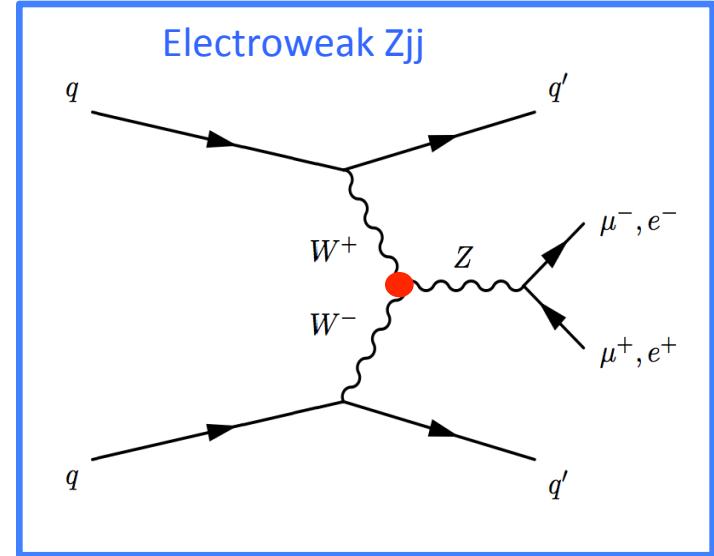
Limits on aTGCs (I)

- Anomalous WWZ couplings parameterised using an effective lagrangian:

$$\frac{\mathcal{L}}{g_{WWZ}} = i \left[g_{1,Z} \left(W_{\mu\nu}^\dagger W^\mu Z^\nu - W_{\mu\nu} W^{\dagger\mu} Z^\nu \right) + \kappa_Z W_\mu^\dagger W_\nu Z^{\mu\nu} + \frac{\lambda_Z}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu Z^{\nu\rho} \right]$$



On-shell W's, off-shell Z
All bosons timelike, $Q^2 > 0$



Off-shell W's, on-shell Z
Z-boson timelike, $Q^2 > 0$
W-bosons space-like, $Q^2 < 0$

Limits on aTGCs (II)

- Number of observed events in data at $m_{jj} > 1$ TeV used to set limits on the aTGC parameters.
 - SHERPA used to parameterise the m_{jj} dependence on the aTGC.
 - Dipole form factor with two choices of unitarisation scale, $\Lambda=6$ TeV and $\Lambda=\infty$
 - Electroweak cross section also measured in this region, for good measure:

aTGC	$\Lambda = 6$ TeV (obs)	$\Lambda = 6$ TeV (exp)	$\Lambda = \infty$ (obs)	$\Lambda = \infty$ (exp)
$\Delta g_{1,Z}$	[-0.65, 0.33]	[-0.58, 0.27]	[-0.50, 0.26]	[-0.45, 0.22]
λ_Z	[-0.22, 0.19]	[-0.19, 0.16]	[-0.15, 0.13]	[-0.14, 0.11]

$$\sigma_{\text{EW}}^{m_{jj}>1\text{TeV}} = 10.7 \pm 0.9 \text{ (stat)} \pm 1.9 \text{ (syst)} \pm 0.3 \text{ (lumi)}$$

Summary

- Inclusive Zjj production measured using ATLAS data at $\sqrt{s}=8\text{TeV}$
 - Cross sections in five fiducial regions
 - Differential distributions fully corrected for detector effects
- Observation of electroweak Zjj production at ATLAS
 - *A benchmark process for future studies of weak-boson fusion at the LHC*
 - Background-only hypothesis rejected at greater than 5σ
 - Cross section measured in two fiducial regions. Excellent agreement with NLO.
 - First limits placed on anomalous triple gauge couplings using weak boson fusion.
- See JHEP 1404 (2014) 031 [arXiv:1401.7610] for more details.